

Therefore,

$$\int (u + u'^2) dx - \int (u_0 + u_0'^2) dx = \int v'^2 dx \geq 0$$

with equality only if $v' \equiv 0$. In this admissible class, this means $v \equiv 0$. \square

1.7 Free boundary conditions

We have considered weak extremals u for which

$$\delta \mathcal{F}_u[\phi] = 0 \quad \text{for all } \phi \in C_c^\infty(a, b).$$

Here we consider the much stronger condition

$$\delta \mathcal{F}_u[\phi] = 0 \quad \text{for all } \phi \in C^1[a, b]. \quad (1.26)$$

Theorem 5 (Proposition 1.12 in BGH). *If $u \in C^1[a, b]$, the Lagrangian F satisfies $F \in C^1([a, b] \times \mathbb{R} \times \mathbb{R})$, and (1.26) holds, then in addition to the Euler-Lagrange equation we have also the endpoints conditions*

$$F_p(a, u(a), u'(a)) = 0 = F_p(b, u(b), u'(b)).$$

Proof: Recall that we have shown by the Lemma of DuBois-Reymond, and the extremality condition, that

$$\psi(x) = F_p(x, u, u') \quad \text{is a } C^1 \text{ function of } x \text{ on } (a, b).$$

In fact,

$$F_p(x, u, u') = F_p(a, u(a), u'(a)) + \int_a^x F_z(\xi, u(\xi), u'(\xi)) d\xi.$$

Notice that ψ has well defined values at the endpoints a and b , and we can calculate right and left derivatives at these points:

$$\psi'(a^+) = \lim_{h \searrow 0} \frac{1}{h} \int_a^{a+h} F_z(\xi, u(\xi), u'(\xi)) d\xi = F_z(a, u(a), u'(a)),$$

and

$$\begin{aligned} \psi'(b^-) &= \lim_{h \searrow 0} \frac{1}{(-h)} \left[\int_a^{b-h} F_z(\xi, u(\xi), u'(\xi)) d\xi - \int_a^b F_z(\xi, u(\xi), u'(\xi)) d\xi \right] \\ &= F_z(b, u(b), u'(b)). \end{aligned}$$

In view of these calculations, we may assume $\psi \in C^1[a, b]$ and return to the extremality condition

$$\int_a^b [F_p(x, u, u')\phi' + F_z(x, u, u')\phi] dx = 0.$$

Since we know extra regularity for $\psi = F_p(x, u, u')$, we can integrate by parts in the usual direction (but this time with $\phi(a)$ and/or $\phi(b)$ possibly nonzero) to get

$$F_p(x, u, u') \Big|_{x=a}^b + \int_a^b \left[F_z(x, u, u') - \frac{d}{dx} F_p(x, u, u') \right] dx = 0 \quad \text{for all } \phi \in C^1[a, b].$$

The integral vanishes because the integrand is (up to a sign) the operator in the Euler-Lagrange equation applied to an extremal. Therefore,

$$F_p(a, u(a), u'(a))\phi(a) = F_p(b, u(b), u'(b))\phi(b) \quad \text{for all } \phi \in C^1[a, b].$$

Taking ϕ with $\phi(a) = 1$ and $\phi(b) = 0$, we obtain

$$F_p(a, u(a), u'(a)) = 0.$$

Similarly, $F_p(b, u(b), u'(b)) = 0$. \square

Exercise 21. *Extend Theorem 5 to the case where*

$$\mathbf{x} : [a, b] \rightarrow \mathbb{R}^n, \quad \mathbf{x} \in C^1([a, b] \rightarrow \mathbb{R}^n),$$

$F \in C^1([a, b] \times \mathbb{R}^n \times \mathbb{R}^n)$, and \mathbf{x} is a weak extremal with free endpoints, meaning

$$\int_a^b \left[\sum_{j=1}^n F_{p_j}(t, \mathbf{x}, \mathbf{x}')\phi_j' + \sum_{j=1}^n F_{z_j}(t, \mathbf{x}, \mathbf{x}')\phi_j \right] dt = 0$$

for all $\vec{\phi} = (\phi_1, \dots, \phi_n) \in C^1([a, b] \rightarrow \mathbb{R}^n)$.

1.8 Technique of the first integral

The following considerations apply to C^2 (classical) extremals of **autonomous** functionals, i.e., ones with Lagrangians $F = F(z, p)$ with no explicit x dependence.